

The evolutions of spinning bodies moving in rotating black hole spacetimes

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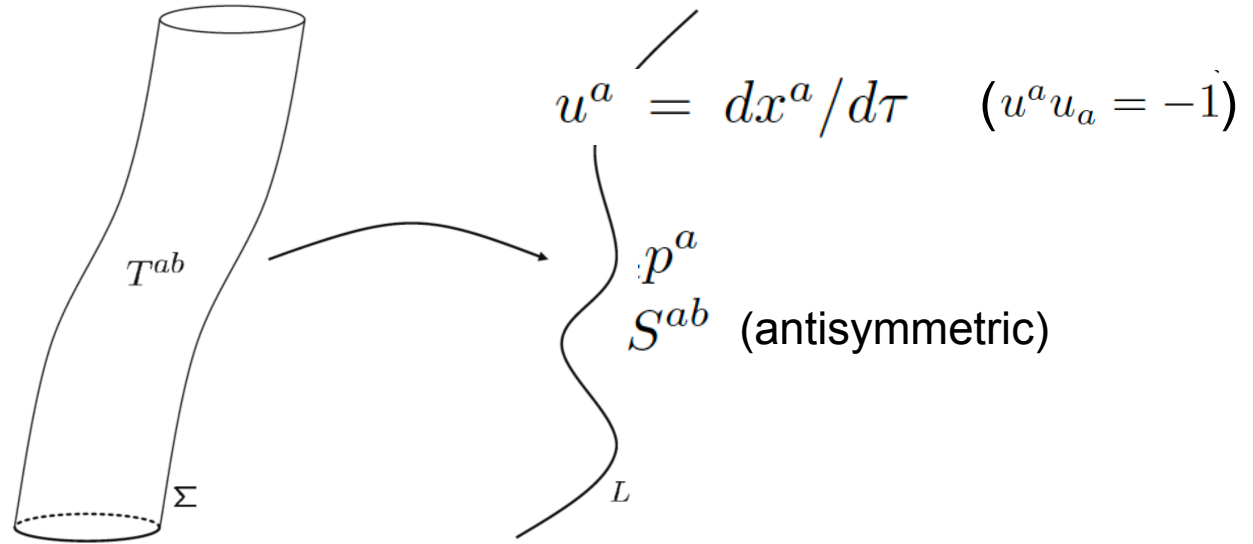
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- The work of Z. K. was supported by the UNKP-18-4 New National Excellence Program of the Ministry of Human Capacities, and by the Janos Bolyai Research Scholarship of the Hungarian Academy of Sciences
- The work of B. M. was supported by the Janos Bolyai Research Scholarship of the Hungarian Academy of Sciences

Second Hermann Minkowski Meeting on the Foundations of Spacetime Physics 2019

Mathisson-Papapetrou-Dixon Eqs.

Description of extended body with multipole moments:



$$\frac{Dp^a}{d\tau} \equiv u^c \nabla_c p^a = F^a, \quad F^a = -\frac{1}{2} R^a{}_{bcd} u^b S^{cd}$$

$$\frac{DS^{ab}}{d\tau} \equiv u^c \nabla_c S^{ab} = p^a u^b - u^a p^b,$$

$$p^a = m u^a - u_b \frac{DS^{ab}}{d\tau}$$

Kinematical mass: $m = -u_a p^a$

Dynamical mass: $M = \sqrt{-p^a p_a}$

M. Mathisson, Acta. Phys. Polon. **6**, 163 (1937).
 A. Papapetrou, Proc. Phys. Soc. **64**, 57 (1951).
 W. Dixon, Nuovo Cim. **34**, 317 (1964).

Spin Supplementary Conditions (SSCs)

Frenkel-Mathisson-Pirani (FMP) SSC

J. Frenkel, Z. Phys. **37**, 243 (1926).
M. Mathisson, Acta. Phys. Polon. **6**, 163 (1937).
F.A.E. Pirani, Acta Phys. Polon. **15**, 389 (1956).

$$u_a S^{ab} = 0$$

Constants: $\frac{D}{d\tau} (S_{ab} S^{ab}) = 0$,

O. Semerák, Mon. Not. Roy. Astron. Soc. **308**, 863 (1999).

$$\frac{dm}{d\tau} = 0$$

Velocity-momentum relation:

L.F.O. Costa, G. Lukes-Gerakopoulos, O. Semerák, Phys. Rev. D **97**, 084023 (2018).

$$m u^a = p^a + \frac{2S^{ab} S_{bc} p^c}{S^{de} S_{de}}$$

Tulczyjew-Dixon (TD) SSC

W.M. Tulczyjew, Acta Phys. Polon. **18**, 393 (1959).
W. Dixon, Nuovo Cim. **34**, 317 (1964).

$$p_a S^{ab} = 0$$

Constants: $\frac{D}{d\tau} (S_{ab} S^{ab}) = 0$,

O. Semerák, Mon. Not. Roy. Astron. Soc. **308**, 863 (1999).

$$\frac{dM}{d\tau} = 0$$

Velocity-momentum relation:

K.P. Tod, F. de Felice, Il Nuovo Cimento **34**, 365 (1976).

$$u^b = \frac{m}{M^2} \left(p^b + \frac{2S^{ba} R_{aecd} p^e S^{cd}}{4M^2 + R_{aecd} S^{ae} S^{cd}} \right)$$

Spin vectors with SSCs

O. Semerák, Mon. Not. Roy. Astron. Soc. **308**, 863 (1999).

Frenkel-Mathisson-Pirani (FMP) SSC

Spin vector:
$$s^a = -\frac{1}{2}\eta^{abcd}u_b S_{cd}$$

$$S^{ab} = \eta^{ab}{}_{cd}u^c s^d$$

Spin magnitude:
$$s_a s^a = \frac{1}{2}S_{cd}S^{cd}$$

Orthogonality relations:

$$s_a S^{ab} = 0, \quad s_a u^a = 0,$$

$$s_a p^a = 0$$

Equation of motion:

$$\frac{Ds^a}{d\tau} = u^a a_b s^b$$

• The case of negligible acceleration was investigated in Ref. D. Bini, A. Geralico, R.T. Jantzen, Phys. Rev. D **95**, 124022 (2017).

Tulczyjew-Dixon (TD) SSC

Spin vector:
$$S^a = -\frac{1}{2M}\eta^{abcd}p_b S_{cd}$$

$$S^{ab} = \frac{1}{M}\eta^{ab}{}_{cd}p^c S^d$$

Spin magnitude:
$$S_a S^a = \frac{1}{2}S_{cd}S^{cd}$$

Orthogonality relations:

$$S_a S^{ab} = 0, \quad p_a S^a = 0,$$

$$S_b u^b = 0$$

Equation of motion:

$$\frac{DS^a}{d\tau} = \frac{S^b F_b}{M^2} p^a$$

Kerr spacetime

R. P. Kerr, Phys. Rev. Lett. 11, 237 (1963).

**Line
element
squared:**

$$ds^2 = -\frac{\Delta - a^2 \sin^2 \theta}{\Sigma} dt^2 - \frac{2a\mathcal{B} \sin^2 \theta}{\Sigma} dt d\phi + \frac{\Sigma}{\Delta} dr^2 + \Sigma d\theta^2 + \frac{\mathcal{A}}{\Sigma} \sin^2 \theta d\phi^2$$

$$\Sigma = r^2 + a^2 \cos^2 \theta, \quad \Delta = r^2 + a^2 - 2\mu r,$$

$$\mathcal{B} = r^2 + a^2 - \Delta, \quad \mathcal{A} = (r^2 + a^2)^2 - \Delta a^2 \sin^2 \theta.$$

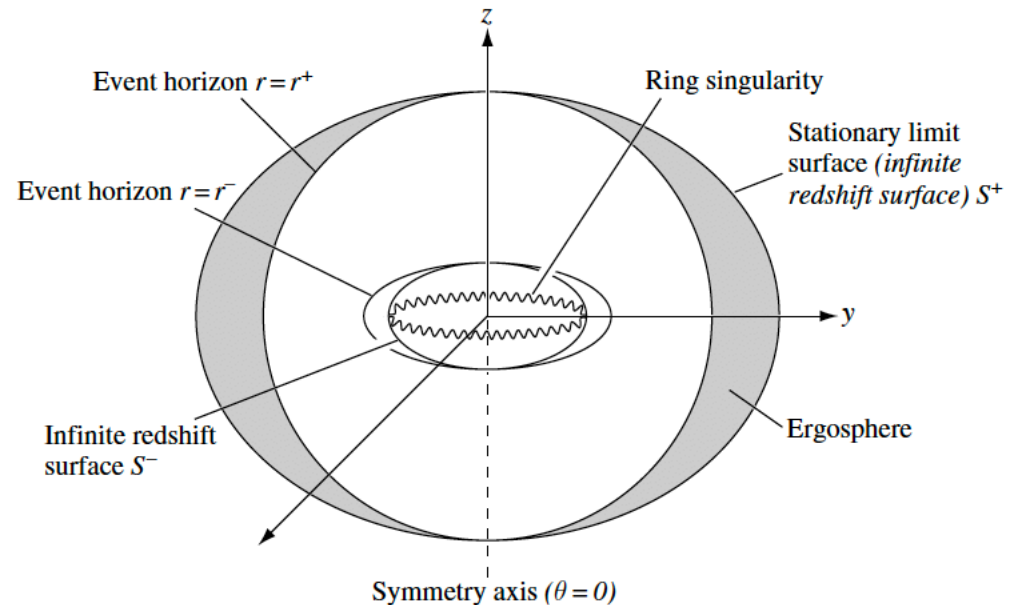
Stationary limit surfaces: $g_{tt} = 0$

Event horizons: $g^{rr} = 0$

Constants of motion:

$$E = -p_t - \frac{1}{2} S^{ab} \partial_a g_{bt},$$

$$J_z = p_\phi + \frac{1}{2} S^{ab} \partial_a g_{b\phi}.$$



Rotating Bardeen-like and Hayward-like spacetimes

B. Toshmatov, Z. Stuchlík, B. Ahmedov, Phys. Rev. D **95**, 084037 (2017).

Line
element
squared:

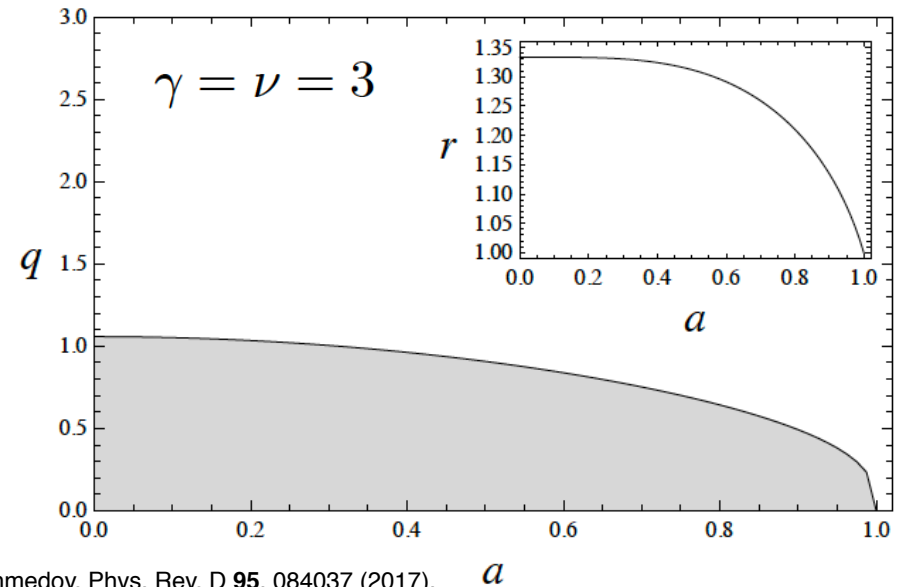
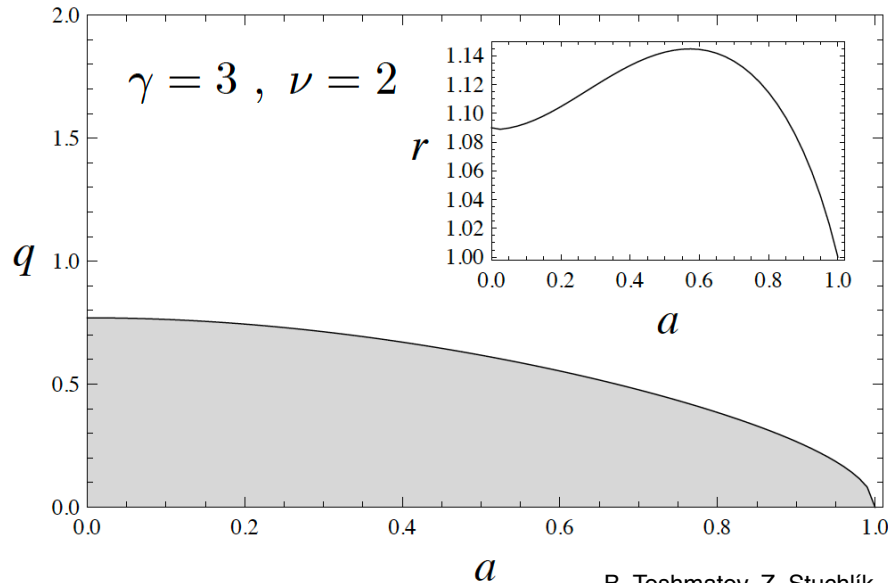
$$ds^2 = -\frac{\Delta - a^2 \sin^2 \theta}{\Sigma} dt^2 - \frac{2a\mathcal{B} \sin^2 \theta}{\Sigma} dt d\phi + \frac{\Sigma}{\Delta} dr^2 + \Sigma d\theta^2 + \frac{\mathcal{A}}{\Sigma} \sin^2 \theta d\phi^2$$

$$\Sigma = r^2 + a^2 \cos^2 \theta, \quad \mathcal{B} = r^2 + a^2 - \Delta, \quad \mathcal{A} = (r^2 + a^2)^2 - \Delta a^2 \sin^2 \theta,$$

$$\Delta = r^2 + a^2 - 2\mu r \frac{r^\gamma}{(r^\nu + q_m^\nu)^{\gamma/\nu}}. \quad \mu = \frac{q_m^3}{\sigma}, \quad q = \frac{q_m}{\mu}$$

Bardeen-like

Hayward-like



B. Toshmatov, Z. Stuchlík, B. Ahmedov, Phys. Rev. D **95**, 084037 (2017).

Comoving and zero 3-momentum frames

Comoving frame: $U^a = u^a$

FMP or TD SSC

zero 3-momentum frame: $U^a = p^a / M$

TD SSC

SO: $u_{(SO)} = \frac{1}{\sqrt{-g_{tt}}} \partial_t$

SO frame vectors:

$$e_0 = u_{(SO)}, \quad e_1 = \sqrt{\frac{\Delta}{\Sigma}} \partial_r, \quad e_2 = \frac{\partial_\theta}{\sqrt{\Sigma}},$$

$$e_3 = -\frac{1}{\sqrt{\Delta}} \left(\frac{a\mathcal{B} \sin \theta}{\Sigma \sqrt{-g_{tt}}} \partial_t - \frac{\sqrt{-g_{tt}}}{\sin \theta} \partial_\phi \right).$$

ZAMO: $u_{(ZAMO)} = \sqrt{\frac{\mathcal{A}}{\Sigma \Delta}} \left(\partial_t + \frac{a\mathcal{B}}{\mathcal{A}} \partial_\phi \right)$

ZAMO frame vectors: f_0, f_α

$$s_a u^a = 0, \quad S_b u^b = 0, \quad p_a S^a = 0$$

Boost transformation: D. Bini, A. Geralico, R.T. Jantzen, Phys. Rev. D **95**, 124022 (2017).

→ $E_0(e, U) \equiv U$

$$E_\alpha(e, U) \quad (\alpha = \{1, 2, 3\})$$

Related by a spatial rotation in U-frame:

Rotation angle: Θ

Boost transformation:

→ $E_0(f, U) \equiv U$

$$E_\alpha(f, U) \quad (\alpha = \{1, 2, 3\})$$

Spin equations in comoving and zero 3-momentum frames

FMP SSC:

$$U^a = u^a$$

$$\frac{ds^\alpha}{d\tau} + \varepsilon^\alpha_{\beta\gamma} \Omega^\beta s^\gamma = 0 .$$

$$\Omega^\alpha = -\frac{1}{2} \varepsilon^{\alpha\beta\gamma} E_\beta \cdot \frac{DE_\gamma}{d\tau}$$

TD SSC:

$$U^a = p^a / M$$

$$\frac{dS^\alpha}{d\tau} + \varepsilon^\alpha_{\beta\gamma} \Omega^\beta S^\gamma = 0 .$$

TD SSC:

$$U^a = u^a \left(\frac{dS^\alpha}{d\tau} + \varepsilon^\alpha_{\beta\gamma} \Omega^\beta S^\gamma \right) E_\alpha + \Upsilon = 0 ,$$

$$\Upsilon = \left[(S \cdot \mathbf{a}) u^{\mathbf{A}} - (S \cdot \mathbf{F}) \frac{p^{\mathbf{A}}}{M^2} \right] E_{\mathbf{A}} .$$

$$u \cdot \Upsilon = 0 \longrightarrow \Upsilon = \Upsilon^\alpha E_\alpha \qquad S \cdot \Upsilon = 0 \longrightarrow \omega \times S = \Upsilon$$

TD SSC:

$$U^a = u^a$$

$$\frac{dS^\alpha}{d\tau} + \varepsilon^\alpha_{\beta\gamma} (\Omega^\beta + \omega^\beta) S^\gamma = 0$$

Cartesian-like triads

SO frame:

$$(e_1, e_2, e_3) = (e_x, e_y, e_z) R_{(e)}$$

$$R_{(e)} \equiv R(\theta, \phi)$$

D. Bini, A. Geralico, R.T. Jantzen,
Phys. Rev. D **95**, 124022 (2017).

$$R(\theta, \phi) = \begin{pmatrix} \sin \theta \cos \phi & \cos \theta \cos \phi & -\sin \phi \\ \sin \theta \sin \phi & \cos \theta \sin \phi & \cos \phi \\ \cos \theta & -\sin \theta & 0 \end{pmatrix}$$

Boosted SO frame:

$$(E_1, E_2, E_3) (e, U)$$

$$= (E_x, E_y, E_z) (e, U) R_{(e)}$$

ZAMO frame:

$$(f_1, f_2, f_3) = (f_x, f_y, f_z) R_{(f)}$$

$$R_{(f)} \equiv R(\theta, \phi)$$

Boosted ZAMO frame:

$$(E_1, E_2, E_3) (f, U)$$

$$= (E_x, E_y, E_z) (f, U) R_{(f)}$$

Evolution equations for Cartesian-like triad components of the spin

FMP SSC: $\mathbb{S}^a = s^a$ **TD SSC:** $\mathbb{S}^a = S^a$ $(p_a \mathbb{S}^a = 0 = u_a \mathbb{S}^a)$

In the U-frame: $\mathbb{S} = \mathbb{S}^\alpha E_\alpha = \mathbb{S}^i E_i$, $\mathbb{S}^0 = 0$.
 $(\alpha = \{1, 2, 3\}$, $\mathbf{i} = \{\mathbf{x}, \mathbf{y}, \mathbf{z}\})$

$$\frac{d\mathbb{S}^i}{d\tau} = -R^i{}_\alpha \varepsilon^\alpha{}_{\beta\gamma} \Omega_{(prec)}^\beta \mathbb{S}^\gamma$$

$$\Omega_{(prec)}^\beta = \Omega_{(p)}^\beta + \epsilon \omega^\beta ,$$

$$\Omega_{(p)}^\beta = -\Omega_{(orb)}^\beta + \Omega^\beta , \quad (R^{-1})^\alpha{}_j \frac{dR^j{}_\beta}{d\tau} = \varepsilon^\alpha{}_{\gamma\beta} \Omega_{(orb)}^\gamma , \quad \Omega^\alpha = -\frac{1}{2} \varepsilon^{\alpha\beta\gamma} E_\beta \cdot \frac{DE_\gamma}{d\tau}$$

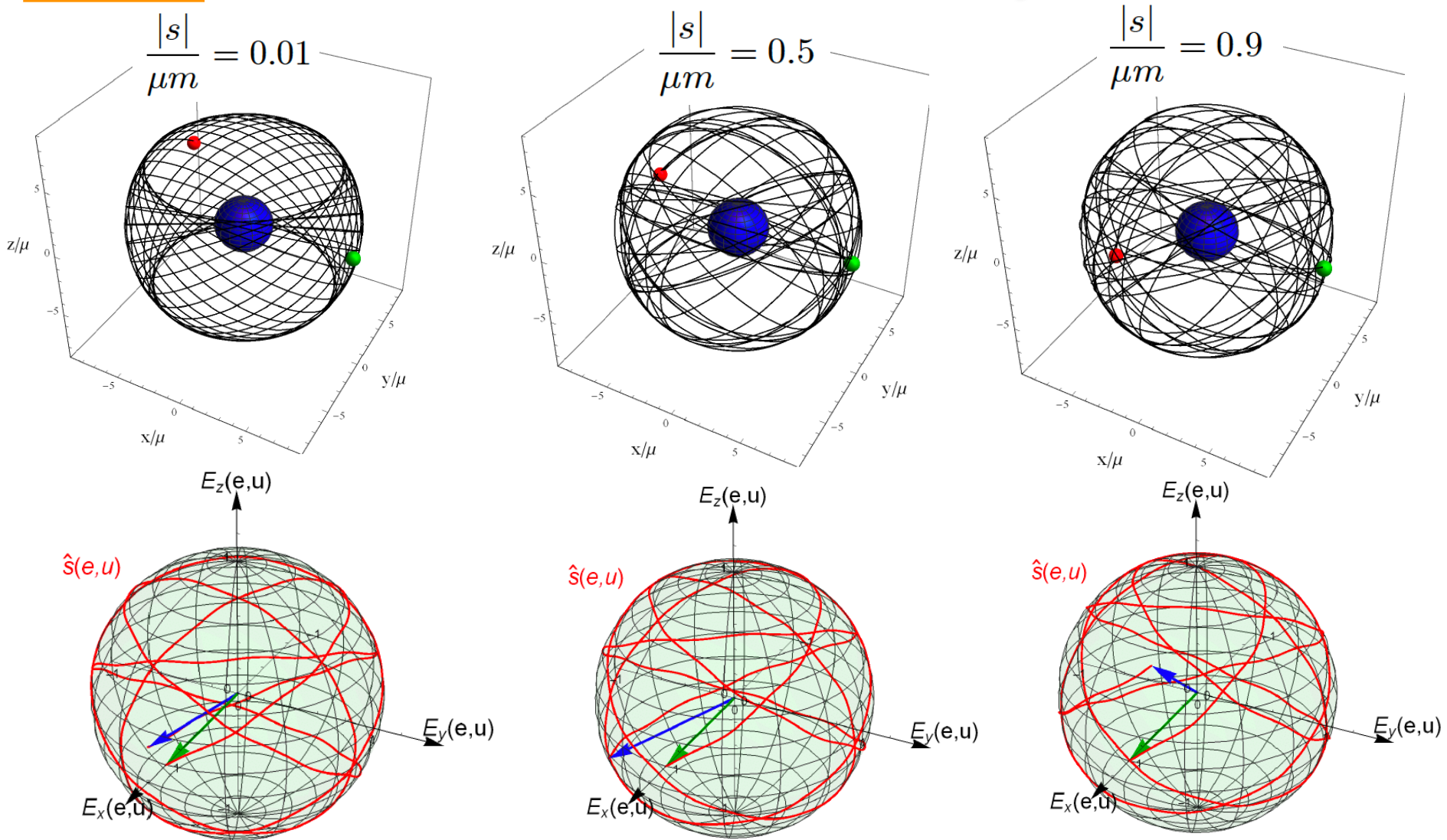
FMP SSC	TD SSC	
$U^a = u^a$	$U^a = p^a / M$	$U^a = u^a$
$\mathbb{S} = s$, $\epsilon = 0$	$\mathbb{S} = S$, $\epsilon = 0$	$\mathbb{S} = S$, $\epsilon = 1$

Spherical-like orbits (Kerr BH)

Coordinate space: $x = r \cos \phi \sin \theta$, $y = r \sin \phi \sin \theta$, $z = r \cos \theta$.

$$a = 0.99\mu$$

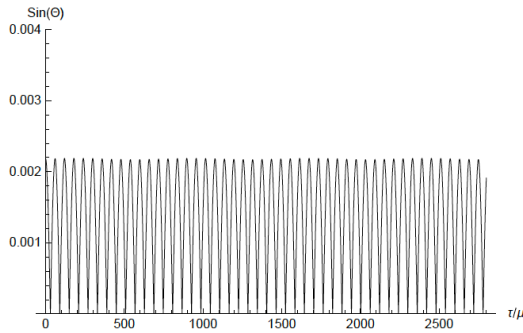
Increasing spin magnitude \longrightarrow



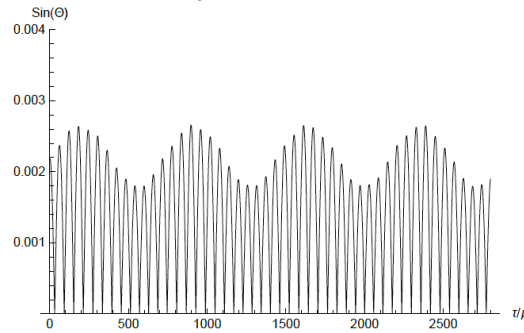
Spherical-like orbits (Kerr BH)

Increasing spin magnitude \longrightarrow

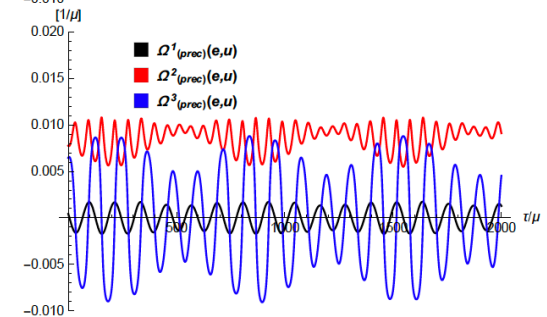
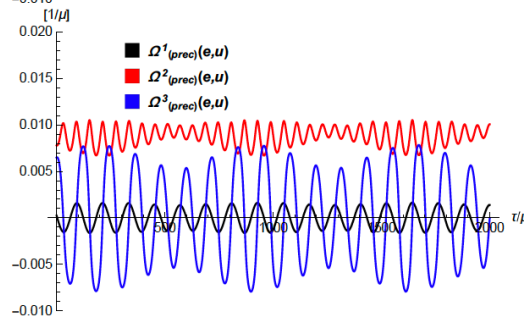
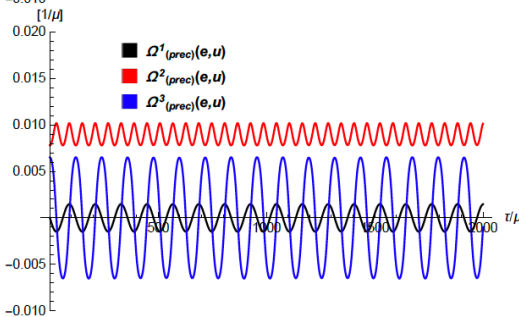
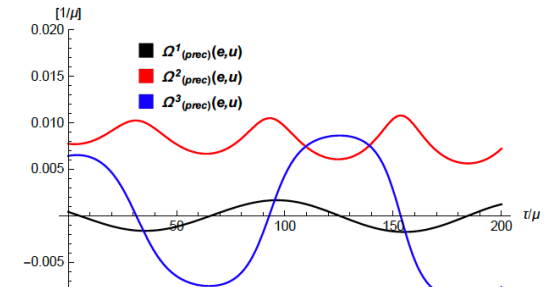
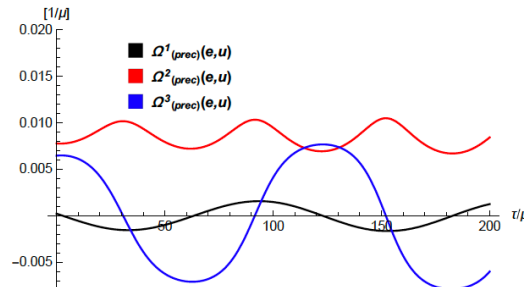
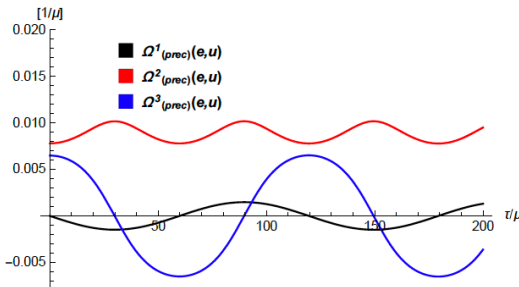
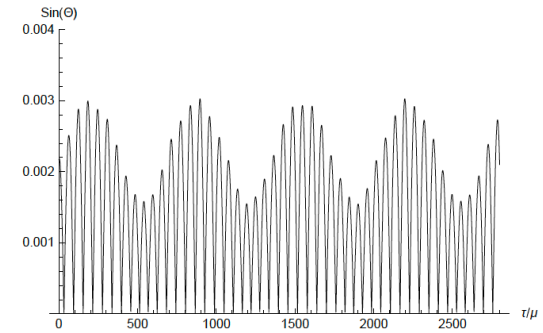
$$\frac{|s|}{\mu m} = 0.01$$



$$\frac{|s|}{\mu m} = 0.5$$



$$\frac{|s|}{\mu m} = 0.9$$

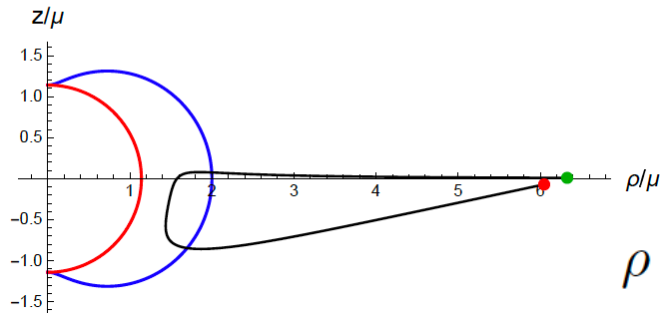
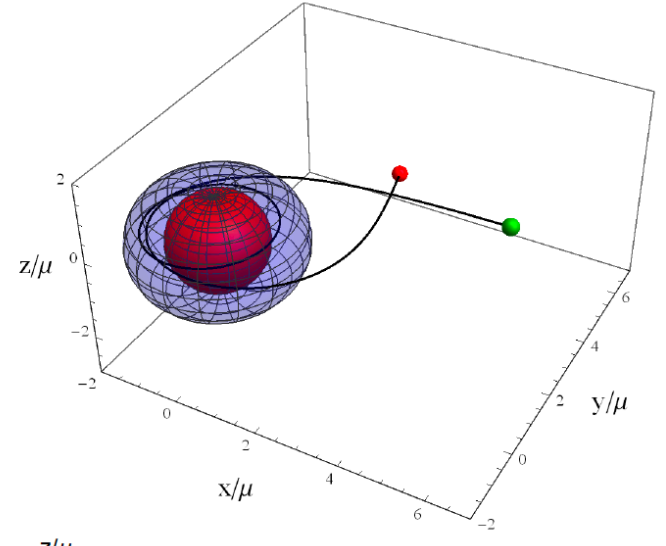
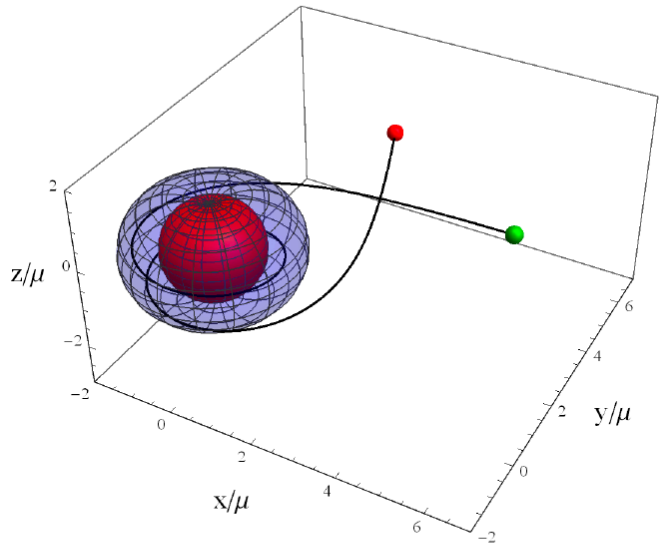
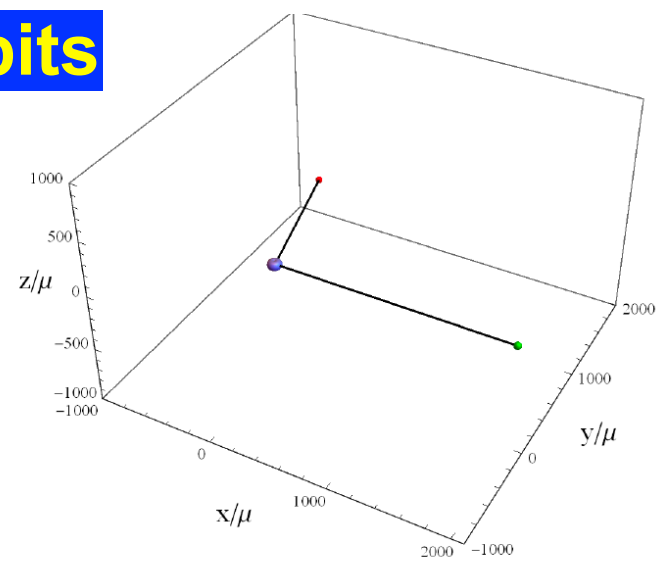
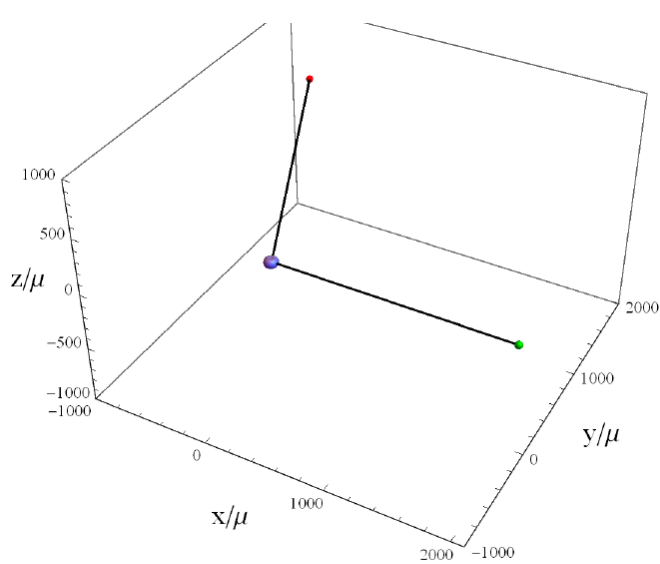


Unbound orbits

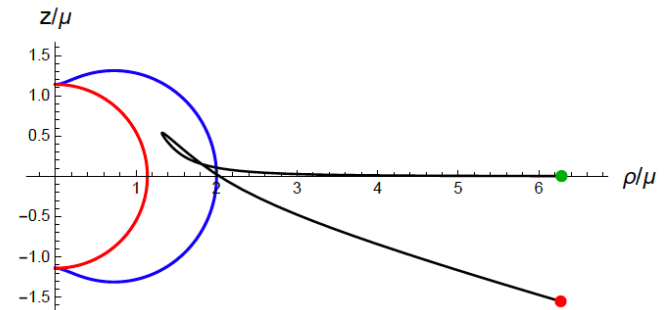
$$a = 0.99\mu$$

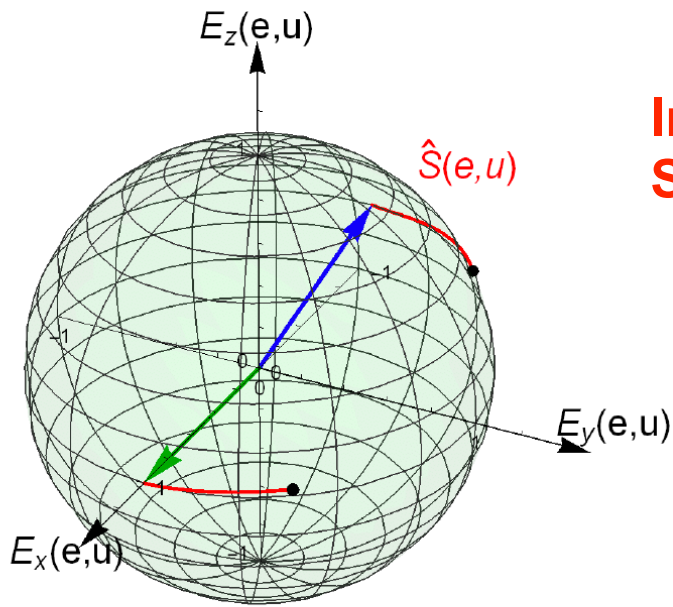
Initial spin directions are different.

$$\frac{|S|}{\mu M} = 0.4$$

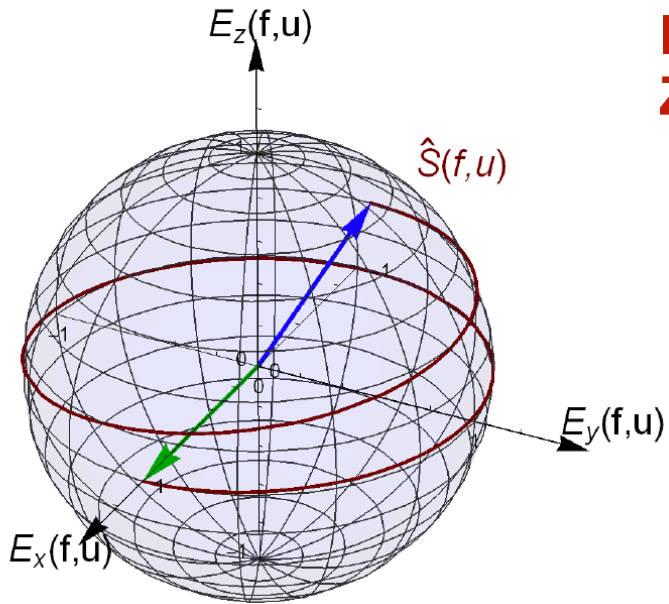
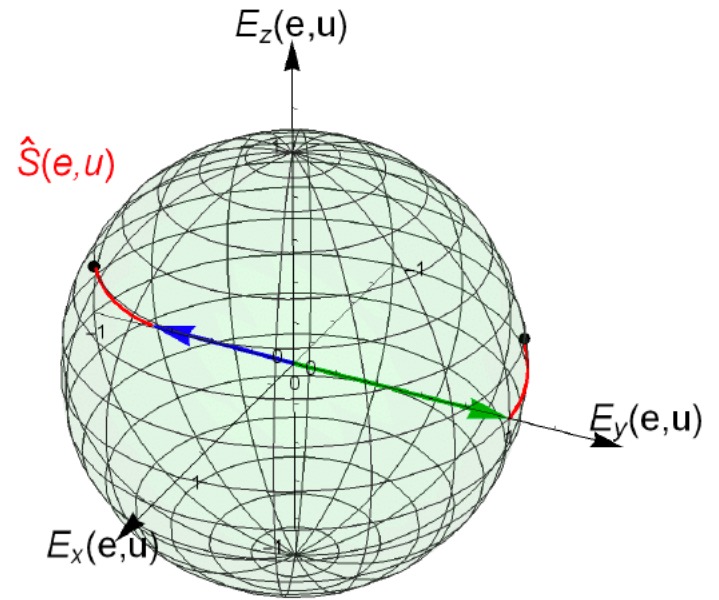


$$\rho = r \sin \theta$$

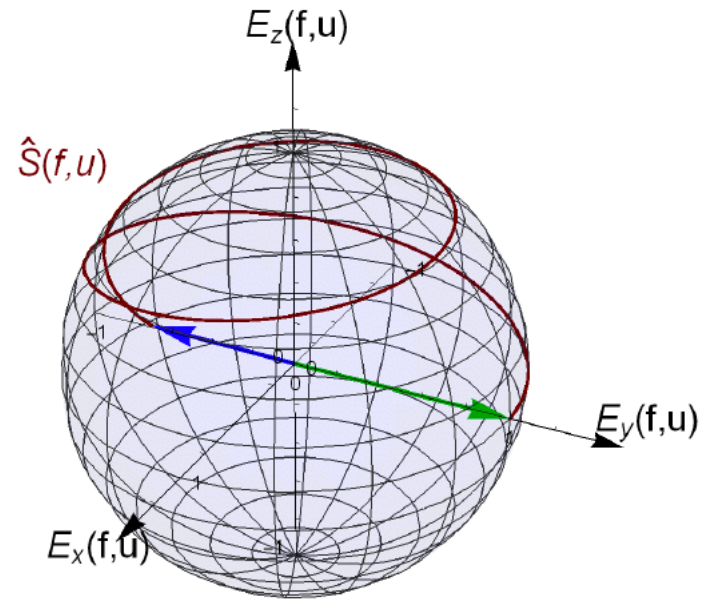




**In boosted
SO frame**

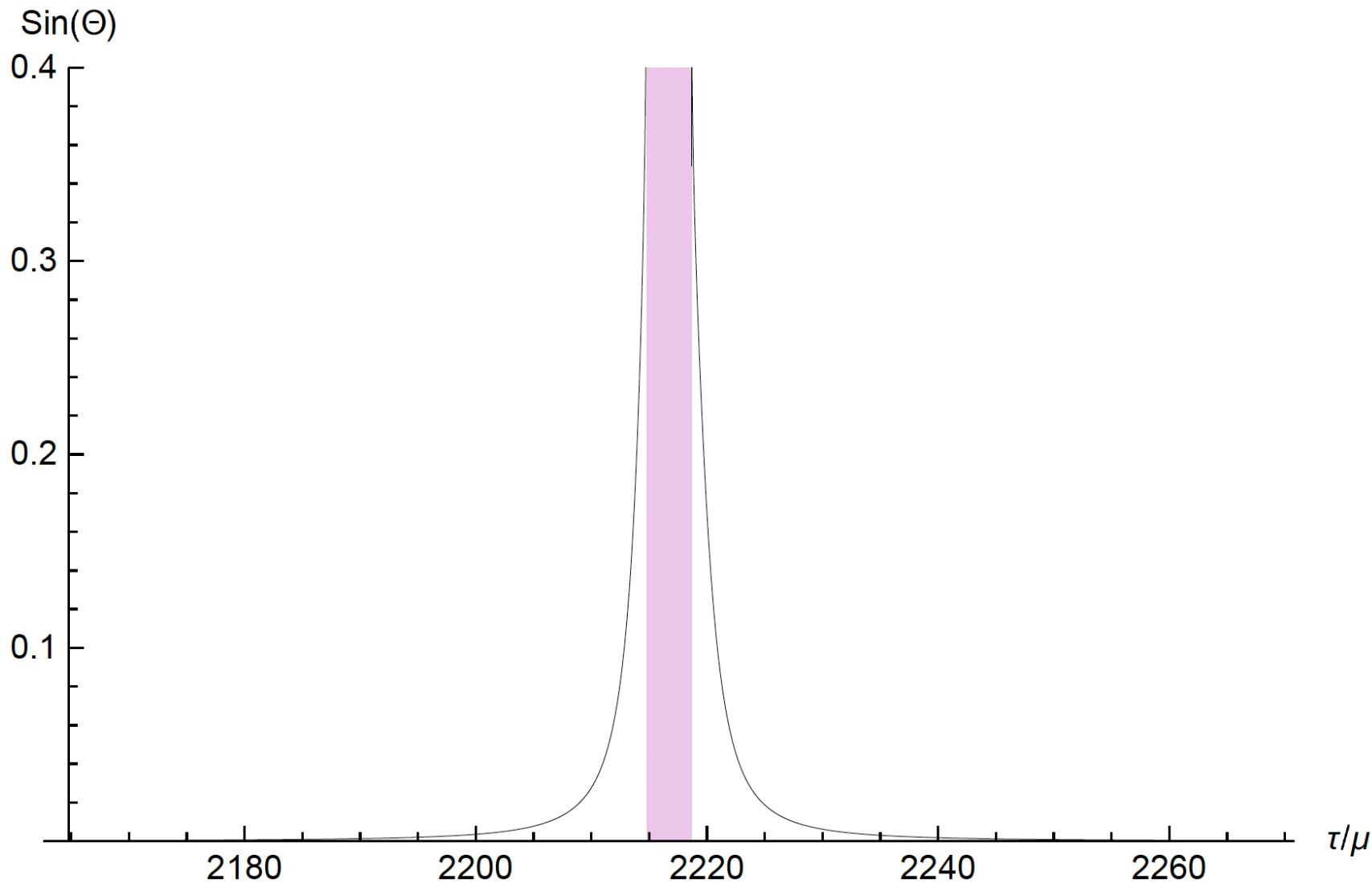


**In boosted
ZAMO frame**



Rotation angle between the boosted SO and ZAMO frames

(In the left side hand case)

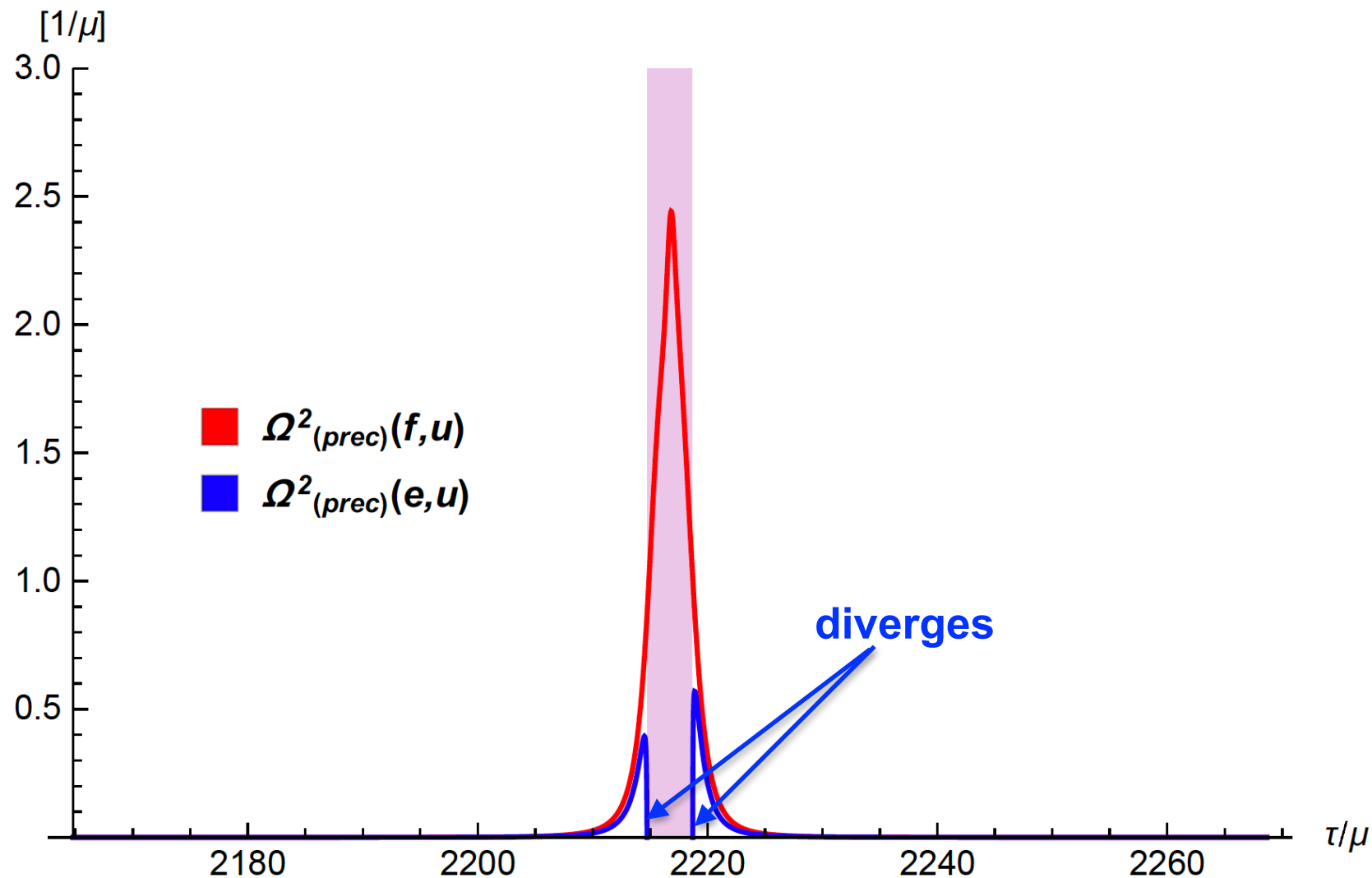


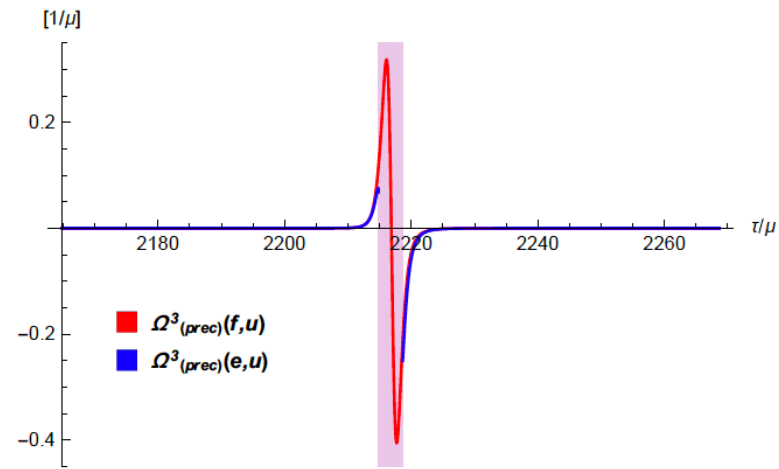
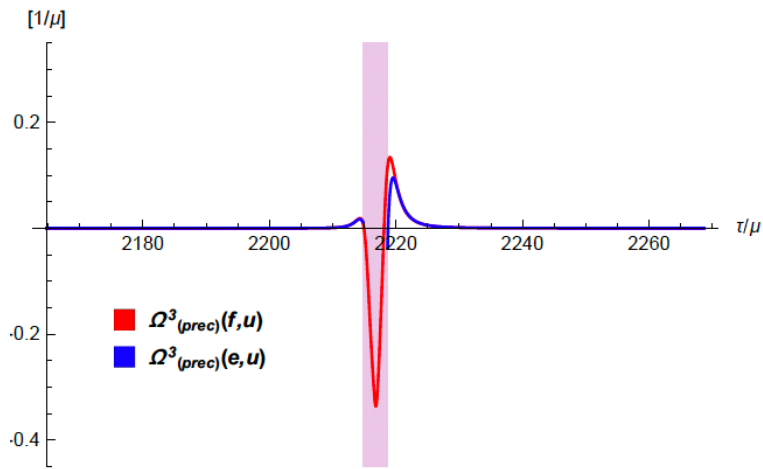
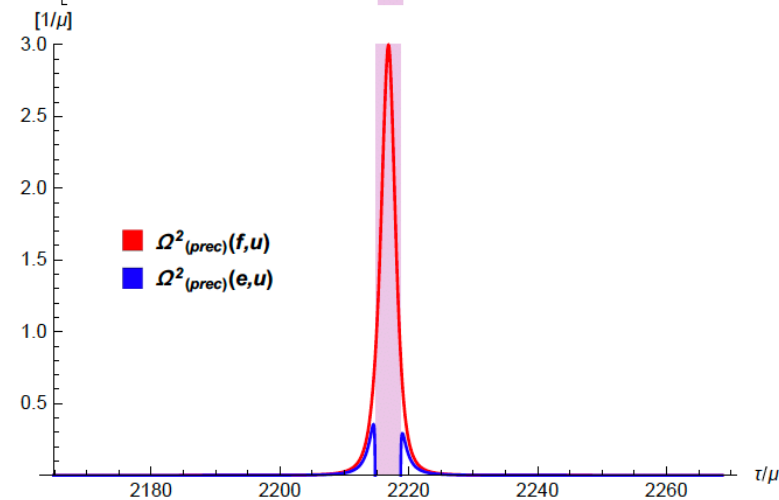
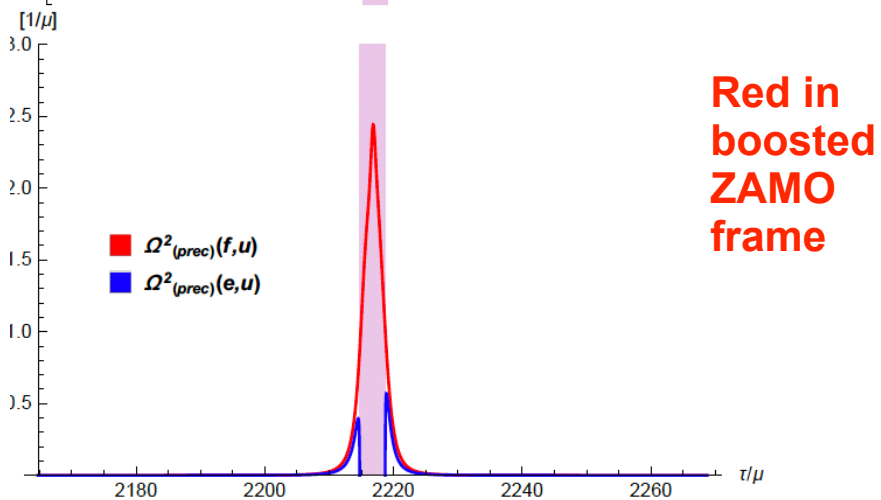
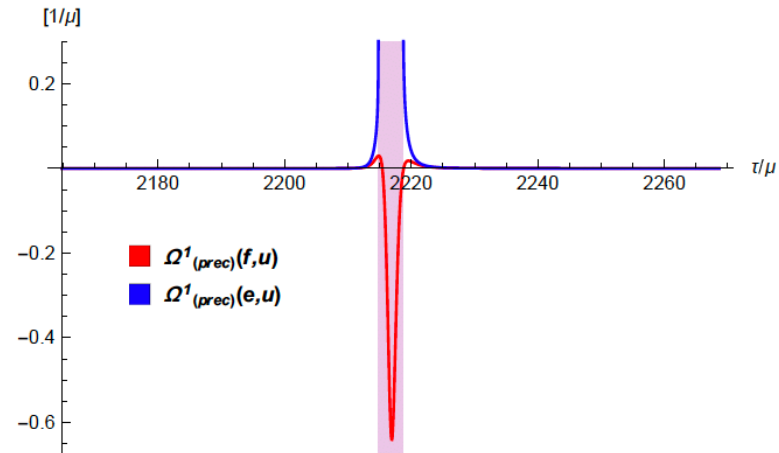
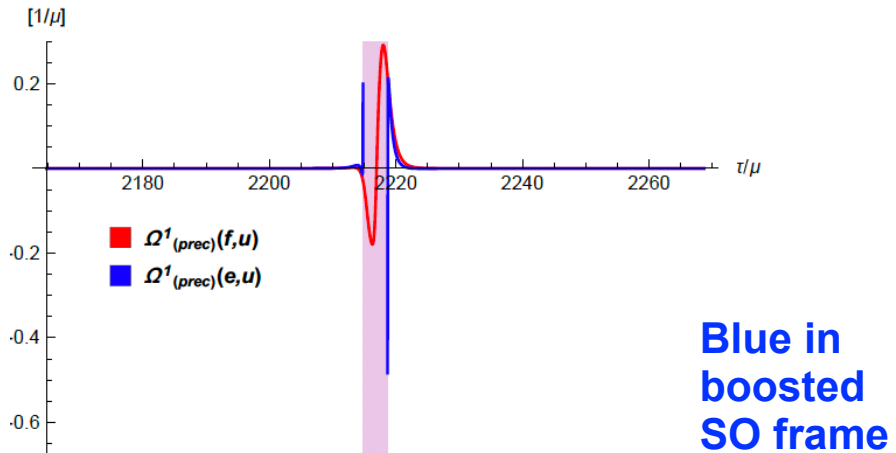
(In the left side hand case)

Precessional angular velocity

Blue in boosted SO frame

Red in boosted ZAMO frame





Zoom-whirl orbits (Kerr and regular BHs)

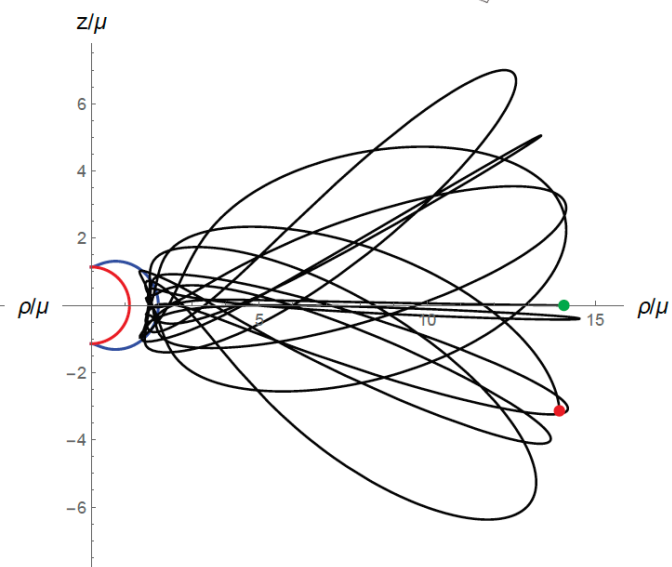
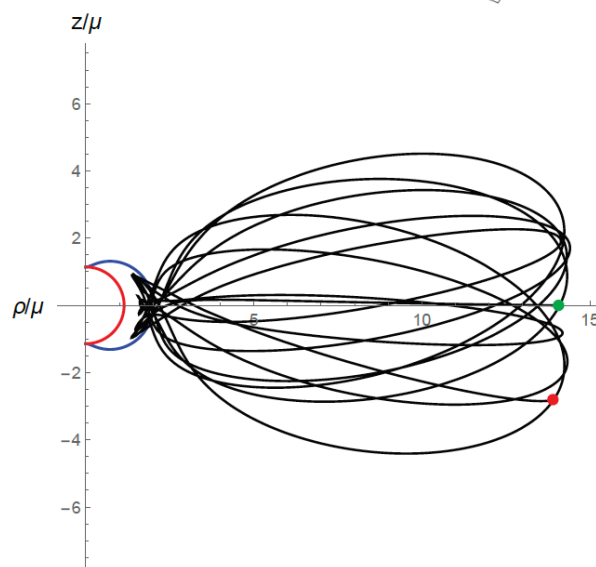
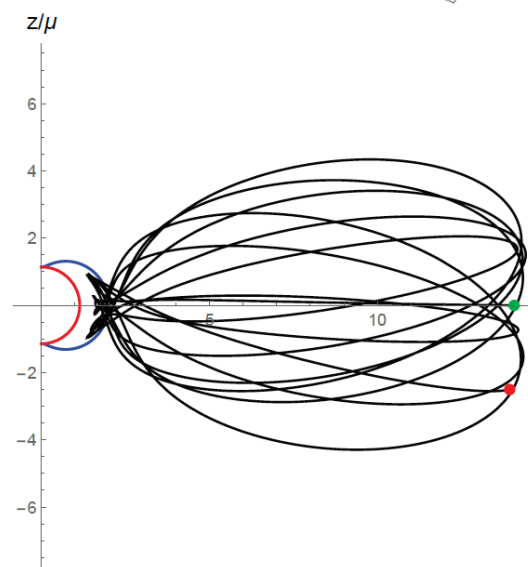
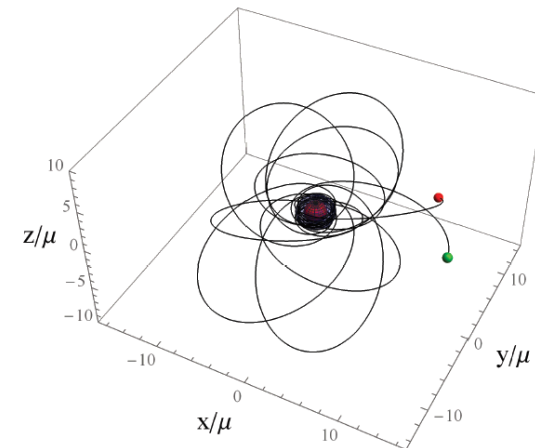
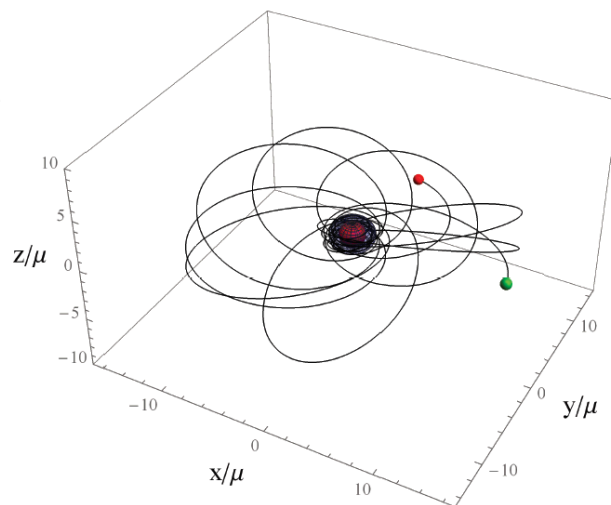
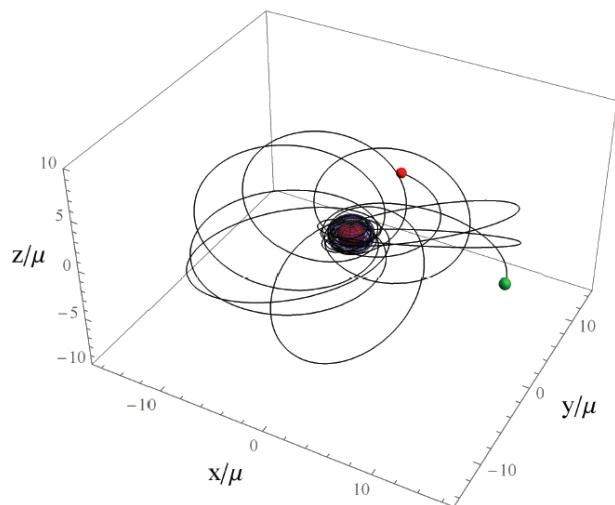
$$a = 0.99\mu$$

$$q = 0.08$$

Kerr BH

Bardeen-like BH

Hayward-like BH

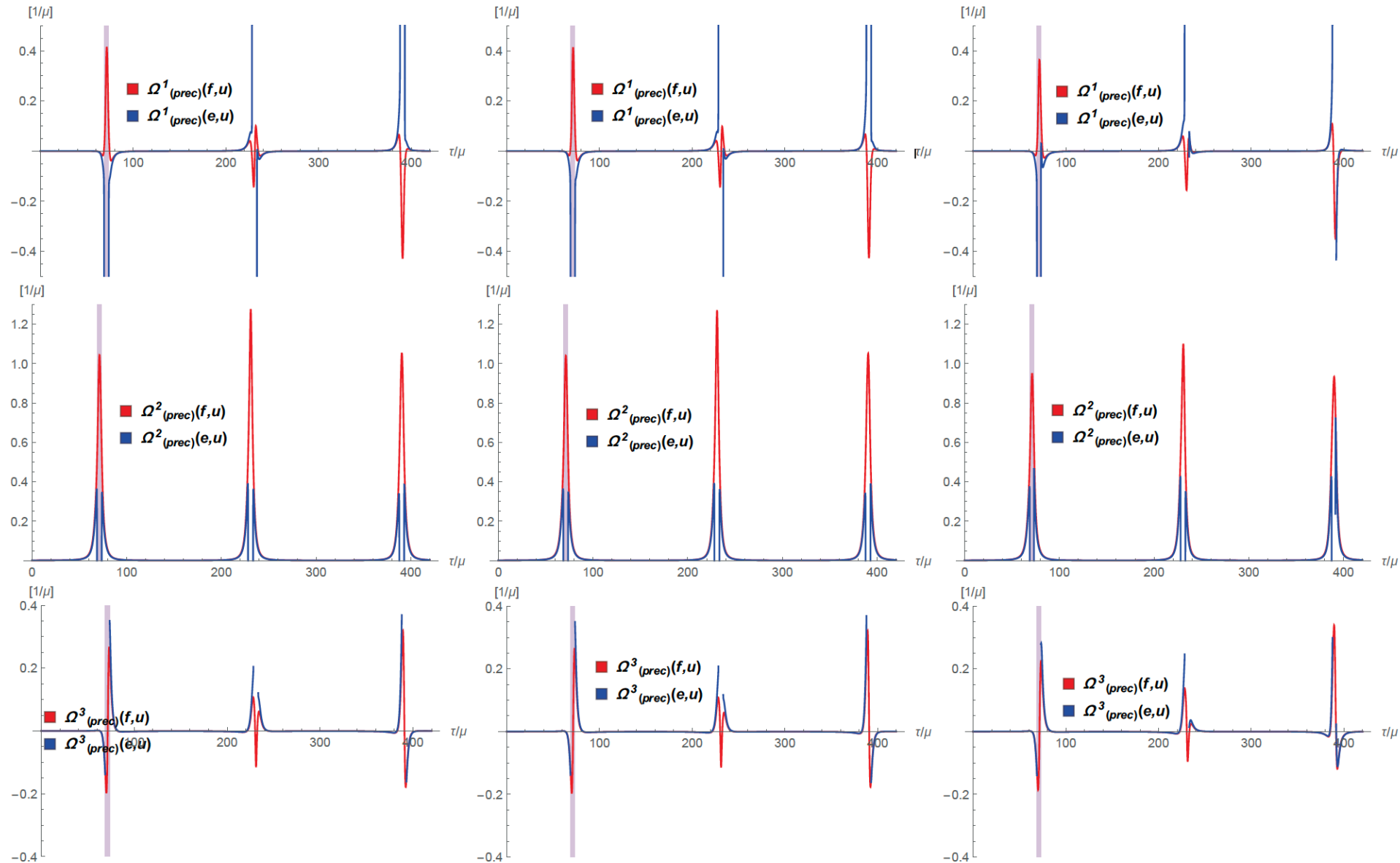


Precessional angular velocity

Kerr BH

Bardeen-like BH

Hayward-like BH





Thank you for the attention

- The work of Z. K. was supported by the UNKP-18-4 New National Excellence Program of the Ministry of Human Capacities, and by the Janos Bolyai Research Scholarship of the Hungarian Academy of Sciences
- The work of B. M. was supported by the Janos Bolyai Research Scholarship of the Hungarian Academy of Sciences