

Thermodynamic manifolds and stability of black holes in various dimensions

T. Vetsov

Department of Physics
Sofia University

“Second Hermann Minkowski Meeting on the Foundations of Spacetime Physics”
Albena, Bulgaria

May 14, 2019

Black hole thermodynamics

Black holes have entropy (Bekenstein-Hawking '70s):

$$S = k_B \frac{A}{4L_p^2} + \text{corrections} \quad (1)$$

The first law of thermodynamics:

$$dM = TdS + \Omega dJ + \Phi dQ + \dots = TdS + \Phi_i dQ^i = I_a dE^a \quad (2)$$

Black holes thermal stability (Davies '80):

$$C = T \frac{\partial S}{\partial T} \begin{cases} > 0, & \text{stable,} \\ < 0, & \text{radiating (unstable),} \\ = 0, & \text{phase transitions,} \\ \rightarrow \infty, & \text{phase transitions} \end{cases} \quad (3)$$

Geometric approaches to the equilibrium space of black holes

The space of extensive parameters $\mathcal{E} = \{\Xi, E^a\}$ is called an equilibrium manifold if supplied with a proper metric structure.

- Hessian information metrics, “Geometric thermodynamics” (F. Weinhold 1975, G. Ruppeiner 1979)
- Legendre invariant metrics, “Geometrothermodynamics” (H. Quevedo 2006)
- Method of conjugate potentials, “New geometric thermodynamics” (B. Mirza, A. Mansoori 2014 & 2019)

Hessian metrics

Fluctuation theory (G. Ruppeiner '79):

$$\begin{aligned} S(E^a) &= S_0 + EQL + \frac{\partial^2 S}{\partial E^a \partial E^b} dE^a dE^b + \dots \\ &= S_0 + EQL - g_{ab}(\vec{E}) dE^a dE^b \end{aligned} \quad (4)$$

Ruppeiner information metric:

$$g_{ab}^{(R)} = -\frac{\partial^2 S}{\partial E^a \partial E^b} = -\text{Hess}S(\vec{E}) \quad (5)$$

Weinhold information metric (F. Weinhold '75):

$$g_{ab}^{(W)} = \frac{\partial^2 M}{\partial E^a \partial E^b} = \text{Hess}M(\vec{E}) \quad (6)$$

Scalar curvature and quantum gravity

- 1 The probability for fluctuating between two macro states is proportional to the geodesic distance between them in \mathcal{E} .
- 2 The strength of interactions/correlations between quantum bits on the event horizon = the magnitude of R
- 3 The sign of R indicates the type of interactions (G. Ruppeiner '10):

$$R \begin{cases} > 0, & \text{repulsive interactions,} \\ < 0, & \text{attractive interactions,} \\ = 0, & \text{free theory,} \\ \rightarrow \infty, & \text{phase transitions} \end{cases} \quad (7)$$

- 4 Phase transitions = divergencies of R (F. Weinhold '75, G. Ruppeiner '79)

Legendre invariant metrics

- Consider $(2n + 1)$ TD phase space \mathcal{T} with coordinates $Z^A = (\Xi, I^a, E^a)$, $a = 1, \dots, n$, where Ξ is a TD potential.
- Select on \mathcal{T} a non-degenerate Legendre invariant metric $G = G(Z^A)$ and a Gibbs 1-form $\Theta(Z^A)$, namely

$$G^{GTD} = \Theta^2 + (\xi_{ab} E^a I^b)(\eta_{cd} dE^c dI^d), \quad \Theta = d\Xi - \delta_{ab} I^a dE^b,$$

where δ_{ab} is the identity matrix, η_{ab} is the Minkowski metric, and ξ_{ab} is some constant tensor.

- Take the pullback $\phi^* : \mathcal{T} \rightarrow \mathcal{E}$ to find (H. Quevedo '17):

$$ds^2 = \left(\delta_{ac} \xi^{cb} E^a \frac{\partial \Xi}{\partial E^b} \right) \left(\eta_e^d \frac{\partial^2 \Xi}{\partial E^d \partial E^f} dE^e dE^f \right) \quad (8)$$

Conjugate thermodynamic potentials

For general black holes with $(m + 1)$ TD variables, (S, Φ_i) , and Enthalpy potential, $\bar{M} = M - \Phi_i Q_i$, one can define the metric (B. Mirza, A. Mansoori '19):

$$\hat{g} = \text{blockdiag} \left(\frac{1}{T} \frac{\partial^2 M}{\partial S^2}, -\hat{G} \right), \quad (9)$$

where

$$G_{ij} = \frac{1}{T} \frac{\partial^2 M}{\partial Y^i \partial Y^j}, \quad Y^i = (Q_1, \dots, Q_m) \quad (10)$$

Black holes in 3 and 4 dimensions

- 1 TIG for 4d Deser-Sarioglu-Tekin black hole solution in higher derivative theory of gravity (T. Vetsov '19):

$$I = \frac{1}{2} \int_{\mathcal{M}} d^4x \sqrt{-g} \left(R + \sigma \sqrt{3 \text{Tr}(\hat{C}^2)} \right), \quad \sigma < -\frac{1}{2} \ \& \ \sigma > 1$$

- 2 TIG for 3d WAdS₃ black hole solution in TMG dual to WCFT₂ with left and right central charges (H. Dimov, R. C. Rashkov, T. Vetsov '19)

$$I = \frac{1}{16\pi G} \int_{\mathcal{M}} d^3x \sqrt{-g} \left(R + \frac{2}{L} \right) + \frac{1}{\mu} I_{CS} + \int_{\partial\mathcal{M}} B$$

$$T_c = \frac{1}{\pi(c_L + \sqrt{c_L c_R})}$$

Summary

- Thermodynamic information geometry (TIG) is a set of geometric tools for investigating statistical thermal systems in equilibrium or non-equilibrium
- TIG is a subset in the more powerful framework of Information Geometry (IG)
- IG is essential for understanding how classical and quantum information can be encoded onto the degrees of freedom of any physical system
- Growing number of applications beyond physics

Thank You!

- In coloboration with R. C. Rashkov and H. Dimov:
 - T. Vetsov, Eur. Phys. J. C (2019) 79: 71
 - H. Dimov, R. C. Rashkov, T. Vetsov, 1902.02433 [hep-th]
- Partially supported by
 - The Bulgarian NSF Grant DM 18/1
 - The Sofia University Grant 10-80-149